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NEW GLOBAL STABILITY ESTIMATES FOR THE CALDERÓN PROBLEM IN TWO DIMENSIONS

MATTEO SANTACESARIA

ABSTRACT. We prove a new global stability estimate for the Gel'fand-Calderón inverse problem on a two-dimensional bounded domain. Specifically, the inverse boundary value problem for the equation $-\Delta\psi + v\psi = 0$ on D is analysed, where v is a smooth real-valued potential of conductivity type defined on a bounded planar domain D . The main feature of this estimate is that it shows that the more a potential is smooth, the more its reconstruction is stable. Furthermore, the stability is proven to depend exponentially on the smoothness, in a sense to be made precise. The same techniques yield a similar estimate for the Calderón problem for the electrical impedance tomography.

1. INTRODUCTION

Let $D \subset \mathbb{R}^2$ be a bounded domain equipped with a potential given by a function $v \in L^\infty(D)$. The corresponding Dirichlet-to-Neumann map is the operator $\Phi : H^{1/2}(\partial D) \rightarrow H^{-1/2}(\partial D)$, defined by

$$(1.1) \quad \Phi(f) = \frac{\partial u}{\partial \nu} \Big|_{\partial D},$$

where $f \in H^{1/2}(\partial D)$, ν is the outer normal of ∂D , and u is the $H^1(D)$ -solution of the Dirichlet problem

$$(1.2) \quad (-\Delta + v)u = 0 \text{ on } D, \quad u|_{\partial D} = f.$$

Here we have assumed that

$$(1.3) \quad 0 \text{ is not a Dirichlet eigenvalue for the operator } -\Delta + v \text{ in } D.$$

The following inverse boundary value problem arises from this construction:

Problem 1. *Given Φ , find v on D .*

This problem can be considered as the Gel'fand inverse boundary value problem for the Schrödinger equation at zero energy (see [10], [17]) as well

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as a generalization of the Calderón problem for the electrical impedance tomography (see [7], [17]), in two dimensions.

It is convenient to recall how the above problem generalises the inverse conductivity problem proposed by Calderón. In the latter, D is a body equipped with an isotropic conductivity $\sigma(x) \in L^\infty(D)$ (with $\sigma \geq \sigma_{\min} > 0$),

$$(1.4) \quad v(x) = \frac{\Delta \sigma^{1/2}(x)}{\sigma^{1/2}(x)}, \quad x \in D,$$

$$(1.5) \quad \Phi = \sigma^{-1/2} \left(\Lambda \sigma^{-1/2} + \frac{\partial \sigma^{1/2}}{\partial \nu} \right),$$

where $\sigma^{-1/2}$, $\partial \sigma^{1/2} / \partial \nu$ in (1.5) denote the multiplication operators by the functions $\sigma^{-1/2}|_{\partial D}$, $\partial \sigma^{1/2} / \partial \nu|_{\partial D}$, respectively and Λ is the voltage-to-current map on ∂D , defined as

$$(1.6) \quad \Lambda f = \sigma \frac{\partial u}{\partial \nu} \Big|_{\partial D},$$

where $f \in H^{1/2}(\partial D)$, ν is the outer normal of ∂D , and u is the $H^1(D)$ -solution of the Dirichlet problem

$$(1.7) \quad \operatorname{div}(\sigma \nabla u) = 0 \text{ on } D, \quad u|_{\partial D} = f.$$

Indeed, the substitution $u = \tilde{u} \sigma^{-1/2}$ in (1.7) yields $(-\Delta + v)\tilde{u} = 0$ in D with v given by (1.4). The following problem is called the Calderón problem:

Problem 2. *Given Λ , find σ on D .*

We remark that Problems 1 and 2 are not overdetermined, in the sense that we consider the reconstruction of a real-valued function of two variables from real-valued inverse problem data dependent on two variables. In addition, the history of inverse problems for the two-dimensional Schrödinger equation at fixed energy goes back to [8].

There are several questions to be answered in these inverse problems: to prove the uniqueness of their solutions (e.g. the injectivity of the map $v \rightarrow \Phi$ for Problem 1), the reconstruction and the stability of the inverse map.

In this paper we study interior stability estimates for the two problems. Let us consider, for instance, Problem 1 with a potential of conductivity type. We want to prove that given two Dirichlet-to-Neumann operators, respectively Φ_1 and Φ_2 , corresponding to potentials, respectively v_1 and v_2 on D , we have that

$$\|v_1 - v_2\|_{L^\infty(D)} \leq \omega(\|\Phi_1 - \Phi_2\|_{H^{1/2} \rightarrow H^{-1/2}}),$$

where the function $\omega(t) \rightarrow 0$ as fast as possible as $t \rightarrow 0$. For Problem 2 similar estimates are considered.

There is a wide literature on the Gel'fand-Calderón inverse problem. In the case of complex-valued potentials the global injectivity of the map $v \rightarrow \Phi$ was firstly proved in [17] for $D \subset \mathbb{R}^d$ with $d \geq 3$ and in [6] for $d = 2$ with $v \in L^p$: in particular, these results were obtained by the use of global reconstructions developed in the same papers. A global stability estimate for Problem 1 and 2 for $d \geq 3$ was first found by Alessandrini in [1]; this result was recently improved in [21]. In the two-dimensional case the first global stability estimate for Problem 1 was given in [23].

Global results for Problem 2 in the two dimensional case have been found much earlier than for Problem 1. In particular, global uniqueness was first proved in [16] for conductivities in the $W^{2,p}(D)$ class ($p > 1$) and after in [2] for L^∞ conductivities. The first global stability result was given in [14], where a logarithmic estimate is obtained for conductivities with two continuous derivatives. This result was improved in [4], where the same kind of estimate is obtained for Hölder continuous conductivities.

The research line delineated above is devoted to prove stability estimates for the least regular potentials/conductivities possible. Here, instead, we focus on the opposite situation, i.e. smooth potentials/conductivities, and try to answer another question: how the stability estimates vary with respect to the smoothness of the potentials/conductivities.

The results, detailed below, also constitute a progress for the case of non-smooth potentials: they indicate stability dependence of the smooth part of a singular potential with respect to boundary value data.

We will assume for simplicity that

$$(1.8) \quad \begin{aligned} &D \text{ is an open bounded domain in } \mathbb{R}^2, \quad \partial D \in C^2, \\ &v \in W^{m,1}(\mathbb{R}^2) \text{ for some } m > 2, \quad \text{supp } v \subset D, \end{aligned}$$

where

$$(1.9) \quad \begin{aligned} W^{m,1}(\mathbb{R}^2) &= \{v : \partial^J v \in L^1(\mathbb{R}^2), |J| \leq m\}, \quad m \in \mathbb{N} \cup \{0\}, \\ J &\in (\mathbb{N} \cup \{0\})^2, \quad |J| = J_1 + J_2, \quad \partial^J v(x) = \frac{\partial^{|J|} v(x)}{\partial x_1^{J_1} \partial x_2^{J_2}}. \end{aligned}$$

Let

$$\|v\|_{m,1} = \max_{|J| \leq m} \|\partial^J v\|_{L^1(\mathbb{R}^2)}.$$

The last (strong) hypothesis is that we will consider only potentials of conductivity type, i.e.

$$(1.10) \quad v = \frac{\Delta \sigma^{1/2}}{\sigma^{1/2}}, \text{ for some } \sigma \in L^\infty(D), \text{ with } \sigma \geq \sigma_{\min} > 0.$$

The main results are the following.

Theorem 1.1. *Let the conditions (1.3), (1.8), (1.10) hold for the potentials v_1, v_2 , where D is fixed, and let Φ_1, Φ_2 be the corresponding Dirichlet-to-Neumann operators. Let $\|v_j\|_{m,1} \leq N$, $j = 1, 2$, for some $N > 0$. Then there exists a constant $C = C(D, N, m)$ such that*

$$(1.11) \quad \|v_2 - v_1\|_{L^\infty(D)} \leq C(\log(3 + \|\Phi_2 - \Phi_1\|^{-1}))^{-\alpha},$$

where $\alpha = m - 2$ and $\|\Phi_2 - \Phi_1\| = \|\Phi_2 - \Phi_1\|_{H^{1/2} \rightarrow H^{-1/2}}$.

Theorem 1.2. *Let σ_1, σ_2 be two isotropic conductivities such that $\Delta(\sigma_j^{1/2})/\sigma_j^{1/2}$ satisfies conditions (1.8), where D is fixed and $0 < \sigma_{\min} \leq \sigma_j \leq \sigma_{\max} < +\infty$ for $j = 1, 2$ and some constants σ_{\min} and σ_{\max} . Let Λ_1, Λ_2 be the corresponding Dirichlet-to-Neumann operators and $\|\Delta(\sigma_j^{1/2})/\sigma_j^{1/2}\|_{m,1} \leq N$, $j = 1, 2$, for some $N > 0$. We suppose, for simplicity, that $\text{supp}(\sigma_j - 1) \subset D$ for $j = 1, 2$. Then, for any $\alpha < m$ there exists a constant $C = C(D, N, \sigma_{\min}, \sigma_{\max}, m, \alpha)$ such that*

$$(1.12) \quad \|\sigma_2 - \sigma_1\|_{L^\infty(D)} \leq C(\log(3 + \|\Lambda_2 - \Lambda_1\|^{-1}))^{-\alpha},$$

where $\|\Lambda_2 - \Lambda_1\| = \|\Lambda_2 - \Lambda_1\|_{H^{1/2} \rightarrow H^{-1/2}}$.

The main feature of these estimates is that, as $m \rightarrow +\infty$, we have $\alpha \rightarrow +\infty$. In addition we would like to mention that, under the assumptions of Theorems 1.1 and 1.2, according to instability estimates of Mandache [15] and Isaev [13], our results are almost optimal. Note that, in the linear approximation near the zero potential, Theorem 1.1 (without condition (1.10)) was proved in [22]. In dimension $d \geq 3$ a global stability estimate similar to our result (with respect to dependence on smoothness) was proved in [21].

The proof of Theorem 1.1 relies on the $\bar{\partial}$ -techniques introduced by Beals–Coifman [5], Henkin–R. Novikov [12], Grinevich–S. Novikov [11] and developed by R. Novikov [17] and Nachman [16] for solving the Calderón problem in two dimensions.

The Novikov–Nachman method starts with the construction of a special family of solutions $\psi(x, \lambda)$ of equation (1.2), which was originally introduced by Faddeev in [9]. These solutions have an exponential behaviour depending on the complex parameter λ and they are constructed via some function $\mu(x, \lambda)$ (see (2.5)). One of the most important property of $\mu(x, \lambda)$ is that it

satisfies a $\bar{\partial}$ -equation with respect to the variable λ (see equation (2.8)), in which appears the so-called Faddeev generalized scattering amplitude $h(\lambda)$ (defined in (2.6)). On the contrary, if one knows $h(\lambda)$ for every $\lambda \in \mathbb{C}$, it is possible to recover $\mu(x, \lambda)$ via this $\bar{\partial}$ -equation. Starting from these arguments we will prove that the map $h(\lambda) \rightarrow \mu(z, \lambda)$ satisfies an Hölder condition, uniformly in the space variable z . This is done in Section 4.

Another part of the method relates the scattering amplitude $h(\lambda)$ to the Dirichlet-to-Neumann operator Φ . In the present paper this is done using the Alessandrini identity (see [1]) and an estimate of $h(\lambda)$ for high values of $|\lambda|$ given in [19]. We find that the map $\Phi \rightarrow h$ has logarithmic stability in some natural norm (Proposition 3.3). This is explained in Section 3.

The final part of the method for the two problems is quite different. For Problem 2, in order to recover $\sigma(x)$ from $\mu(x, \lambda)$, we use a limit found for the first time in [16]. Instead, for Problem 1, we use an explicit formula for $v(x)$ which involves the scattering amplitude $h(\lambda)$, $\mu(x, \lambda)$ and its first (complex) derivative with respect to $z = x_1 + ix_2$ (see formula (5.3)). The two results are presented in section 5 and yield the proofs of Theorems 1.1 and 1.2.

This work was fulfilled in the framework of researches under the direction of R. G. Novikov.

2. PRELIMINARIES

In this section we recall some definitions and properties of the Faddeev functions, the above-mentioned family of solutions of equation (1.2), which will be used throughout all the paper.

Following [16], we fix some $1 < p < 2$ and define $\psi(x, k)$ to be the solution (when it exists unique) of

$$(2.1) \quad (-\Delta + v)\psi(x, k) = 0 \text{ in } \mathbb{R}^2,$$

with $e^{-ixk}\psi(x, k) - 1 \in W^{1, \bar{p}}(\mathbb{R}^2) = \{u : \partial^J u \in L^{\bar{p}}(\mathbb{R}^2), |J| \leq 1\}$, where $x = (x_1, x_2) \in \mathbb{R}^2$, $k = (k_1, k_2) \in \mathcal{V} \subset \mathbb{C}^2$,

$$(2.2) \quad \mathcal{V} = \{k \in \mathbb{C}^2 : k^2 = k_1^2 + k_2^2 = 0\}$$

and

$$(2.3) \quad \frac{1}{\bar{p}} = \frac{1}{p} - \frac{1}{2}.$$

The variety \mathcal{V} can be written as $\{(\lambda, i\lambda) : \lambda \in \mathbb{C}\} \cup \{(\lambda, -i\lambda) : \lambda \in \mathbb{C}\}$. We henceforth denote $\psi(x, (\lambda, i\lambda))$ by $\psi(x, \lambda)$ and observe that, since v is real-valued, uniqueness for (2.1) yields $\psi(x, (-\bar{\lambda}, i\bar{\lambda})) = \overline{\psi(x, (\lambda, i\lambda))} = \overline{\psi(x, \lambda)}$ so that, for reconstruction and stability purpose, it is sufficient to work on the sheet $k = (\lambda, i\lambda)$.

We now identify \mathbb{R}^2 with \mathbb{C} and use the coordinates $z = x_1 + ix_2$, $\bar{z} = x_1 - ix_2$,

$$\frac{\partial}{\partial z} = \frac{1}{2} \left(\frac{\partial}{\partial x_1} - i \frac{\partial}{\partial x_2} \right), \quad \frac{\partial}{\partial \bar{z}} = \frac{1}{2} \left(\frac{\partial}{\partial x_1} + i \frac{\partial}{\partial x_2} \right),$$

where $(x_1, x_2) \in \mathbb{R}^2$.

Then we define

$$(2.4) \quad \psi(z, \lambda) = \psi(x, \lambda),$$

$$(2.5) \quad \mu(z, \lambda) = e^{-iz\lambda} \psi(z, \lambda),$$

$$(2.6) \quad h(\lambda) = \int_D e^{i\bar{z}\bar{\lambda}} v(z) \psi(z, \lambda) d\operatorname{Re} z d\operatorname{Im} z,$$

for $z, \lambda \in \mathbb{C}$.

Throughout all the paper $c(\alpha, \beta, \dots)$ is a positive constant depending on parameters α, β, \dots

We now restate some fundamental results about Faddeev functions. In the following statement ψ_0 denotes $\sigma^{1/2}$.

Proposition 2.1 (see [16]). *Let $D \subset \mathbb{R}^2$ be an open bounded domain with C^2 boundary, $v \in L^p(\mathbb{R}^2)$, $1 < p < 2$, $\operatorname{supp} v \subset D$, $\|v\|_{L^p(\mathbb{R}^2)} \leq N$, be such that there exists a real-valued $\psi_0 \in L^\infty(\mathbb{R}^2)$ with $v = (\Delta \psi_0)/\psi_0$, $\psi_0(x) \geq c_0 > 0$ and $\psi_0 \equiv 1$ outside D . Then, for any $\lambda \in \mathbb{C}$ there is a unique solution $\psi(z, \lambda)$ of (2.1) with $e^{-iz\lambda} \psi(\cdot, \lambda) - 1$ in $L^{\tilde{p}} \cap L^\infty$ (\tilde{p} is defined in (2.3)). Furthermore, $e^{-iz\lambda} \psi(\cdot, \lambda) - 1 \in W^{1, \tilde{p}}(\mathbb{R}^2)$ and*

$$(2.7) \quad \|e^{-iz\lambda} \psi(\cdot, \lambda) - 1\|_{W^{s, \tilde{p}}} \leq c(p, s) N |\lambda|^{s-1},$$

for $0 \leq s \leq 1$ and λ sufficiently large.

The function $\mu(z, \lambda)$ defined in (2.5) satisfies the equation

$$(2.8) \quad \frac{\partial \mu(z, \lambda)}{\partial \bar{\lambda}} = \frac{1}{4\pi \bar{\lambda}} h(\lambda) e_{-\lambda}(z) \overline{\mu(z, \lambda)}, \quad z, \lambda \in \mathbb{C},$$

in the $W^{1, \tilde{p}}$ topology, where $h(\lambda)$ is defined in (2.6) and the function $e_{-\lambda}(z)$ is defined as follows:

$$(2.9) \quad e_\lambda(z) = e^{i(z\lambda + \bar{z}\bar{\lambda})}.$$

In addition, the functions $h(\lambda)$ and $\mu(z, \lambda)$ satisfy

$$(2.10) \quad \left\| \frac{h(\lambda)}{\bar{\lambda}} \right\|_{L^r(\mathbb{R}^2)} \leq c(r, N), \text{ for all } r \in (\tilde{p}', \tilde{p}), \quad \frac{1}{\tilde{p}} + \frac{1}{\tilde{p}'} = 1,$$

$$(2.11) \quad \sup_{z \in \mathbb{C}} \|\mu(z, \cdot) - 1\|_{L^r(\mathbb{C})} \leq c(r, D, N), \quad \text{for all } r \in (p', \infty]$$

and

$$(2.12) \quad |h(\lambda)| \leq c(p, D, N)|\lambda|^\varepsilon,$$

$$(2.13) \quad \|\mu(\cdot, \lambda) - \psi_0\|_{W^{1,\bar{p}}} \leq c(p, D, N)|\lambda|^\varepsilon,$$

for $\lambda \leq \lambda_0(p, D, N)$ and $0 < \varepsilon < \frac{2}{p'}$, where $\frac{1}{p} + \frac{1}{p'} = 1$.

Remark. Equation (2.8) means that μ is a generalised analytic function in $\lambda \in \mathbb{C}$ (see [24]). In two-dimensional inverse scattering for the Schrödinger equation, the theory of generalised analytic functions was used for the first time in [11].

We recall that if $v \in W^{m,1}(\mathbb{R}^2)$ with $\text{supp } v \subset D$, then $\|\hat{v}\|_m < +\infty$, where

$$(2.14) \quad \hat{v}(p) = (2\pi)^{-2} \int_{\mathbb{R}^2} e^{ipx} v(x) dx, \quad p \in \mathbb{C}^2,$$

$$(2.15) \quad \|u\|_m = \sup_{p \in \mathbb{R}^2} |(1 + |p|^2)^{m/2} u(p)|,$$

for a test function u .

In addition, if $v \in W^{m,1}(\mathbb{R}^2)$ with $\text{supp } v \subset D$ and $m > 2$, we have, by Sobolev embedding, that

$$(2.16) \quad \|v\|_{L^\infty(D)} \leq c(D) \|v\|_{m,1},$$

so, in particular, the hypothesis $v \in L^p(\mathbb{R}^2)$, $\text{supp } v \subset D$, in the statement of Proposition 2.1 is satisfied for every $1 < p < 2$ (since D is bounded).

The following lemma is a variation of a result in [19]:

Lemma 2.2. *Under the assumption (1.8), there exists $R = R(m, \|\hat{v}\|_m) > 0$ such that*

$$(2.17) \quad |h(\lambda)| \leq 8\pi^2 \|\hat{v}\|_m (1 + 4|\lambda|^2)^{-m/2}, \quad \text{for } |\lambda| > R.$$

Proof. We consider the function $H(k, p)$ defined as

$$(2.18) \quad H(k, p) = \frac{1}{(2\pi)^2} \int_{\mathbb{R}^2} e^{i(p-k)x} v(x) \psi(x, k) dx,$$

for $k \in \mathcal{V}$ (where \mathcal{V} is defined in (2.2)), $p \in \mathbb{R}^2$ and $\psi(x, k)$ as defined at the beginning of this section.

We deduce that $h(\lambda) = (2\pi)^2 H(k(\lambda), k(\lambda) + \overline{k(\lambda)})$, for $k(\lambda) = (\lambda, i\lambda)$. By [19, Corollary 1.1] we have

$$(2.19) \quad |H(k, p)| \leq 2\|\hat{v}\|_m (1 + p^2)^{-m/2} \quad \text{for } |\lambda| > R,$$

for $R = R(m, \|\hat{v}\|_m) > 0$ and then the proof follows. \square

We restate [3, Lemma 2.6], which will be useful in section 4.

Lemma 2.3 ([3]). *Let $a \in L^{s_1}(\mathbb{R}^2) \cap L^{s_2}(\mathbb{R}^2)$, $1 < s_1 < 2 < s_2 < \infty$ and $b \in L^s(\mathbb{R}^2)$, $1 < s < 2$. Assume u is a function in $L^{\tilde{s}}(\mathbb{R}^2)$, with \tilde{s} defined as in (2.3), which satisfies*

$$(2.20) \quad \frac{\partial u(\lambda)}{\partial \lambda} = a(\lambda)\bar{u}(\lambda) + b(\lambda), \quad \lambda \in \mathbb{C}.$$

Then there exists $c > 0$ such that

$$(2.21) \quad \|u\|_{L^{\tilde{s}}} \leq c\|b\|_{L^s} \exp(c(\|a\|_{L^{s_1}} + \|a\|_{L^{s_2}})).$$

We will make also use of the well-known Hölder's inequality, which we recall in a special case: for $f \in L^p(\mathbb{C})$, $g \in L^q(\mathbb{C})$ such that $1 \leq p, q \leq \infty$, $1 \leq r < \infty$, $1/p + 1/q = 1/r$, we have

$$\|fg\|_{L^r(\mathbb{C})} \leq \|f\|_{L^p(\mathbb{C})} \|g\|_{L^q(\mathbb{C})}.$$

3. FROM Φ TO $h(\lambda)$

Lemma 3.1. *Let the condition (1.8) holds. Then we have, for $p \geq 1$,*

$$(3.1) \quad \left\| \frac{h(\lambda)}{\lambda} \right\|_{L^p(|\lambda| > R)} \leq c(p, m) \|\hat{v}\|_m \frac{1}{R^{m+1-2/p}},$$

$$(3.2) \quad \|h\|_{L^p(|\lambda| > R)} \leq c(p, m) \|\hat{v}\|_m \frac{1}{R^{m-2/p}},$$

where R is as in Lemma 2.2.

Proof. It's a corollary of Lemma 2.2. Indeed we have

$$(3.3) \quad \left\| \frac{h(\lambda)}{\lambda} \right\|_{L^p(|\lambda| > R)}^p \leq c \|\hat{v}\|_m^p \int_{r>R} r^{1-mp-p} dr = \frac{c(p, m) \|\hat{v}\|_m^p}{R^{(m+1)p-2}},$$

which gives (3.1). The proof of (3.2) is analogous. \square

Lemma 3.2. *Let $D \subset \{x \in \mathbb{R}^2 : |x| \leq l\}$, v_1, v_2 be two potentials satisfying (1.3), (1.8), (1.10), let Φ_1, Φ_2 the corresponding Dirichlet-to-Neumann operator and h_1, h_2 the corresponding generalised scattering amplitude. Let $\|v_j\|_{m,1} \leq N$, $j = 1, 2$. Then we have*

$$(3.4) \quad |h_2(\lambda) - h_1(\lambda)| \leq c(D, N) e^{2l|\lambda|} \|\Phi_2 - \Phi_1\|_{H^{1/2} \rightarrow H^{-1/2}}, \quad \lambda \in \mathbb{C}.$$

Proof. We have the following identity:

$$(3.5) \quad h_2(\lambda) - h_1(\lambda) = \int_{\partial D} \overline{\psi_1(z, \lambda)} (\Phi_2 - \Phi_1) \psi_2(z, \lambda) |dz|,$$

where $\psi_j(z, \lambda)$ are the Faddeev functions associated to the potential v_j , $j = 1, 2$. This identity is a particular case of the one in [20, Theorem 1]: we refer to that paper for a proof.

From this identity we have:

$$(3.6) \quad |h_2(\lambda) - h_1(\lambda)| \leq \|\psi_1(\cdot, \lambda)\|_{H^{1/2}(\partial D)} \|\Phi_2 - \Phi_1\|_{H^{1/2} \rightarrow H^{-1/2}} \|\psi_2(\cdot, \lambda)\|_{H^{1/2}(\partial D)}.$$

Now take $\tilde{p} > 2$ and use the trace theorem to get

$$\begin{aligned} \|\psi_j(\cdot, \lambda)\|_{H^{1/2}(\partial D)} &\leq C \|\psi_j(\cdot, \lambda)\|_{W^{1, \tilde{p}}(D)} \leq C e^{l|\lambda|} \|e^{-iz\lambda} \psi_j(\cdot, \lambda)\|_{W^{1, \tilde{p}}(D)} \\ &\leq C e^{l|\lambda|} \left(\|e^{-iz\lambda} \psi_j(\cdot, \lambda) - 1\|_{W^{1, \tilde{p}}(D)} + \|1\|_{W^{1, \tilde{p}}(D)} \right), \quad j = 1, 2, \end{aligned}$$

which from (2.7) and (2.11) is bounded by $C(D, N) e^{l|\lambda|}$. These estimates together with (3.6) give (3.4). \square

The main results of this section are the following propositions:

Proposition 3.3. *Let v_1, v_2 be two potentials satisfying (1.3), (1.8), (1.10), let Φ_1, Φ_2 the corresponding Dirichlet-to-Neumann operator and h_1, h_2 the corresponding generalised scattering amplitude. Let $0 < \varepsilon < 1$, $1 < p < \frac{2}{1-\varepsilon}$ and $\|v_j\|_{m,1} \leq N$, $j = 1, 2$. Then there exists a constant $c = c(D, N, m, p)$ such that*

$$(3.7) \quad \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^p(\mathbb{C})} \leq c \log(3 + \|\Phi_2 - \Phi_1\|_{H^{1/2} \rightarrow H^{-1/2}}^{-1})^{-(m+1-2/p)}.$$

Proposition 3.4. *Let $v_1, v_2, \Phi_1, \Phi_2, h_1, h_2$ be as in Proposition 3.3. Let $p \geq 1$ and $\|v_j\|_{m,1} \leq N$, $j = 1, 2$. Then there exists a constant $c = c(D, N, m, p)$ such that*

$$(3.8) \quad \|h_2 - h_1\|_{L^p(\mathbb{C})} \leq c \log(3 + \|\Phi_2 - \Phi_1\|_{H^{1/2} \rightarrow H^{-1/2}}^{-1})^{-(m-2/p)}.$$

Proof of Proposition 3.3. Let choose $a, b > 0$, a close to 0 and b big to be determined and let

$$(3.9) \quad \delta = \|\Phi_2 - \Phi_1\|_{H^{1/2} \rightarrow H^{-1/2}}.$$

We split down the left term of (3.7) as follows:

$$\begin{aligned} \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^p(\mathbb{C})} &\leq \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^p(|\lambda| < a)} + \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^p(a < |\lambda| < b)} \\ &\quad + \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^p(|\lambda| > b)}. \end{aligned}$$

From (2.12) we obtain

$$\begin{aligned} (3.10) \quad \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^p(|\lambda| < a)} &\leq c(D, N, p) \left(\int_{|\lambda| < a} |\lambda|^{(\varepsilon-1)p} d\operatorname{Re} \lambda d\operatorname{Im} \lambda \right)^{\frac{1}{p}} \\ &= c(D, N, p) a^{\varepsilon-1+2/p}. \end{aligned}$$

From Lemma 3.2 and (3.9) we get

$$(3.11) \quad \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^p(a < |\lambda| < b)} \leq c(D, N) \left(\frac{\delta}{a^{1-2/p}} + \delta e^{2lb} \right).$$

From Lemma 3.1

$$(3.12) \quad \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^p(|\lambda| > b)} \leq \frac{c(N)}{b^{m+1-2/p}}.$$

We now define

$$(3.13) \quad a = \log(3 + \delta^{-1})^{-\frac{m+1-2/p}{\varepsilon-1+2/p}}, \quad b = \beta \log(3 + \delta^{-1}),$$

for $0 < \beta < 1/(2l)$, in order to have (3.10) and (3.12) of the order $\log(3 + \delta^{-1})^{-(m+1-2/p)}$. We also choose $\bar{\delta} < 1$ such that for every $\delta \leq \bar{\delta}$, a is sufficiently small in order to have (2.12) (which yields (3.10)), $b \geq R$ (with R as in Lemma 2.2) and also

$$(3.14) \quad \frac{\delta}{a^{1-2/p}} = \delta \log(3 + \delta^{-1})^{\left(\frac{m+1-2/p}{\varepsilon-1+2/p}\right)(1-2/p)} < \log(3 + \delta^{-1})^{-(m+1-2/p)}.$$

Thus we obtain

$$(3.15) \quad \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^p(\mathbb{C})} \leq \frac{c(D, N, p)}{\log(3 + \delta^{-1})^{m+1-2/p}} + c(D, N) \delta (3 + \delta^{-1})^{2l\beta},$$

for $\delta \leq \bar{\delta}$, $0 < \beta < 1/(2l)$. As $\delta(3 + \delta^{-1})^{2l\beta} \rightarrow 0$ for $\delta \rightarrow 0$ more rapidly than the other term, we obtain that

$$(3.16) \quad \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^p(\mathbb{C})} \leq \frac{c(D, N, m, p, \beta)}{\log(3 + \delta^{-1})^{m+1-2/p}},$$

for $\delta \leq \bar{\delta}$, $0 < \beta < 1/(2l)$.

Estimate (3.16) for general δ (with modified constant) follows from (3.16) for $\delta \leq \bar{\delta}$ and the property (2.10) of the scattering amplitude. This completes the proof of Proposition 3.3. \square

Proof of Proposition 3.4. We follow almost the same scheme as in the proof of Proposition 3.3. Let choose $b > 0$ big to be determined and let

$$(3.17) \quad \delta = \|\Phi_2 - \Phi_1\|_{H^{1/2} \rightarrow H^{-1/2}}.$$

We split down the left term of (3.8) as follows:

$$\|h_2 - h_1\|_{L^p(\mathbb{C})} \leq \|h_2 - h_1\|_{L^p(|\lambda| < b)} + \|h_2 - h_1\|_{L^p(|\lambda| \geq b)}.$$

From Lemma 3.2 we obtain

$$(3.18) \quad \|h_2 - h_1\|_{L^p(|\lambda| < b)} \leq c(D, N, p) \delta b^{1/p} e^{2lb},$$

and from (3.2)

$$(3.19) \quad \|h_2 - h_1\|_{L^p(|\lambda| \geq b)} \leq c(N, p, m) \frac{1}{b^{m-2/p}}.$$

Define $b = \beta \log(3 + \delta^{-1})$ for $0 < \beta < 1/(2l)$. Let $\bar{\delta} < 1$ such that for $\delta \leq \bar{\delta}$ we have that $b > R$, where R is defined in Lemma 2.2.

Then we have, for $\delta \leq \bar{\delta}$,

$$\begin{aligned} \|h_2 - h_1\|_{L^p(\mathbb{C})} &\leq c(D, N, m, p) \delta (1 + \delta^{-1})^{2l\beta} (\beta \log(3 + \delta^{-1}))^{1/p} \\ &\quad + c(N, m, p) (\log(3 + \delta^{-1}))^{-(m-2/p)}. \end{aligned}$$

Since $2l\beta < 1$, we have that

$$\delta (1 + \delta^{-1})^{2l\beta} (\beta \log(3 + \delta^{-1}))^{1/p} \rightarrow 0 \quad \text{for } \delta \rightarrow 0$$

more rapidly than the other term. Thus

$$(3.20) \quad \|h_2 - h_1\|_{L^p(\mathbb{C})} \leq c(D, N, m, p, \beta) (\log(3 + \delta^{-1}))^{-(m-2/p)},$$

for $\delta \leq \bar{\delta}$, $0 < \beta < 1/(2l)$.

Estimate (3.20) for general δ (with modified constant) follows from (3.20) for $\delta \leq \bar{\delta}$ and the L^p -boundedness of the scattering amplitude (this because it is continuous and decays at infinity like in Lemma 3.1). This completes the proof of Proposition 3.4. \square

4. ESTIMATES ON THE FADDEEV FUNCTIONS

Lemma 4.1. *Let v_1, v_2 be two potentials satisfying (1.3), (1.8), (1.10), with $\|v_j\|_{m,1} \leq N$, h_1, h_2 the corresponding scattering amplitude and $\mu_1(z, \lambda), \mu_2(z, \lambda)$ the corresponding Faddeev functions. Let $1 < s < 2$, and \tilde{s} be as in (2.3). Then*

$$(4.1) \quad \sup_{z \in \mathbb{C}} \|\mu_2(z, \cdot) - \mu_1(z, \cdot)\|_{L^{\tilde{s}}(\mathbb{C})} \leq c(D, N, s) \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^s(\mathbb{C})},$$

$$(4.2) \quad \sup_{z \in \mathbb{C}} \left\| \frac{\partial \mu_2(z, \cdot)}{\partial z} - \frac{\partial \mu_1(z, \cdot)}{\partial z} \right\|_{L^{\tilde{s}}(\mathbb{C})} \leq c(D, N, s) \left[\left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^s(\mathbb{C})} + \|h_2 - h_1\|_{L^s(\mathbb{C})} \right]$$

Proof. We begin with the proof of (4.1). Let

$$(4.3) \quad \nu(z, \lambda) = \mu_2(z, \lambda) - \mu_1(z, \lambda).$$

From the $\bar{\partial}$ -equation (2.8) we deduce that ν satisfies the following non-homogeneous $\bar{\partial}$ -equation:

$$(4.4) \quad \frac{\partial}{\partial \bar{\lambda}} \nu(z, \lambda) = \frac{e_{-\lambda}(z)}{4\pi} \left(\frac{h_1(\lambda)}{\bar{\lambda}} \overline{\nu(z, \lambda)} + \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \overline{\mu_2(z, \lambda)} \right),$$

for $\lambda \in \mathbb{C}$, where $e_{-\lambda}(z)$ is defined in (2.9). Note that since, by Sobolev embedding, $v \in L^\infty(D) \subset L^s(D)$, we have that $\nu(z, \cdot) \in L^{\tilde{s}}(\mathbb{C})$ for every $\tilde{s} > 2$ (see (2.11)). In addition, from Proposition 2.1 (see (2.10)) we have that $h(\lambda)/\bar{\lambda} \in L^p(\mathbb{C})$, for $1 < p < \infty$. Then it is possible to use Lemma 2.3 in order to obtain

$$\begin{aligned} \|\nu(z, \cdot)\|_{L^{\tilde{s}}} &\leq c(D, N, s) \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^s(\mathbb{C})} \\ &\leq c(D, N, s) \sup_{z \in \mathbb{C}} \|\mu_2(z, \cdot)\|_{L^\infty} \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^s(\mathbb{C})} \\ &\leq c(D, N, s) \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^s(\mathbb{C})}, \end{aligned}$$

where we used again the property (2.11) of $\mu_2(z, \lambda)$.

Now we pass to (4.2). To simplify notations we write, for $z, \lambda \in \mathbb{C}$,

$$\mu_z^j(z, \lambda) = \frac{\partial \mu_j(z, \lambda)}{\partial z}, \quad \mu_{\bar{z}}^j(z, \lambda) = \frac{\partial \mu_j(z, \lambda)}{\partial \bar{z}}, \quad j = 1, 2.$$

From the $\bar{\partial}$ -equation (2.8) we have that μ_z^j and $\mu_{\bar{z}}^j$ satisfy the following system of non-homogeneous $\bar{\partial}$ -equations, for $j = 1, 2$:

$$\begin{aligned} \frac{\partial}{\partial \bar{\lambda}} \mu_z^j(z, \lambda) &= \frac{e_{-\lambda}(z)}{4\pi} \frac{h_j(\lambda)}{\bar{\lambda}} \left(\overline{\mu_{\bar{z}}^j(z, \lambda)} - i\lambda \overline{\mu_j(z, \lambda)} \right), \\ \frac{\partial}{\partial \bar{\lambda}} \mu_{\bar{z}}^j(z, \lambda) &= \frac{e_{-\lambda}(z)}{4\pi} \frac{h_j(\lambda)}{\bar{\lambda}} \left(\overline{\mu_z^j(z, \lambda)} - i\bar{\lambda} \overline{\mu_j(z, \lambda)} \right). \end{aligned}$$

Define now $\mu_\pm^j(z, \lambda) = \mu_z^j(z, \lambda) \pm \mu_{\bar{z}}^j(z, \lambda)$, for $j = 1, 2$. Then they satisfy the following two non-homogeneous $\bar{\partial}$ -equations:

$$\frac{\partial}{\partial \bar{\lambda}} \mu_\pm^j(z, \lambda) = \pm \frac{e_{-\lambda}(z)}{4\pi} \frac{h_j(\lambda)}{\bar{\lambda}} \left(\overline{\mu_\pm^j(z, \lambda)} \mp i(\lambda \pm \bar{\lambda}) \overline{\mu_j(z, \lambda)} \right).$$

Finally define $\tau_\pm(z, \lambda) = \mu_\pm^2(z, \lambda) - \mu_\pm^1(z, \lambda)$. They satisfy the two non-homogeneous $\bar{\partial}$ -equations below:

$$\begin{aligned} \frac{\partial}{\partial \bar{\lambda}} \tau_\pm(z, \lambda) &= \pm \frac{e_{-\lambda}(z)}{4\pi} \left[\frac{h_1(\lambda)}{\bar{\lambda}} \tau_\pm(z, \lambda) + \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \overline{\mu_\pm^2(z, \lambda)} \right. \\ &\quad \left. \mp i \frac{\lambda \pm \bar{\lambda}}{\bar{\lambda}} \left((h_2(\lambda) - h_1(\lambda)) \overline{\mu_2(z, \lambda)} + h_1(\lambda) \overline{\nu(z, \lambda)} \right) \right], \end{aligned}$$

where $\nu(z, \lambda)$ was defined in (4.3).

Now remark that by [19, Lemma 2.1] and regularity assumptions on the potentials we have that $\mu_z^j(z, \cdot), \mu_{\bar{z}}^j(z, \cdot) \in L^{\tilde{s}}(\mathbb{C}) \cap L^\infty(\mathbb{C})$ for any $\tilde{s} > 2$, $j = 1, 2$. This, in particular, yields $\tau_\pm(z, \cdot) \in L^{\tilde{s}}(\mathbb{C})$. These arguments, along with the above remarks on the L^p boundedness of $h_j(\lambda)/\bar{\lambda}$, make possible to use Lemma 2.3, which gives

$$\begin{aligned} \|\tau_\pm(z, \cdot)\|_{L^{\tilde{s}}(\mathbb{C})} &\leq c(D, N, s) \left[\left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \overline{\mu_\pm^2(z, \cdot)} \right\|_{L^s(\mathbb{C})} \right. \\ &\quad \left. + \|(h_2(\cdot) - h_1(\cdot)) \overline{\mu_2(z, \cdot)}\|_{L^s(\mathbb{C})} + \|h_1(\cdot) \overline{\nu(z, \cdot)}\|_{L^s(\mathbb{C})} \right] \\ &\leq c(D, N, s) \left[\left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^s(\mathbb{C})} + \|h_2 - h_1\|_{L^s(\mathbb{C})} \right. \\ &\quad \left. + \|h_1\|_{L^2(\mathbb{C})} \|\nu(z, \cdot)\|_{L^{\tilde{s}}(\mathbb{C})} \right] \\ &\leq c(D, N, s) \left[\left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^s(\mathbb{C})} + \|h_2 - h_1\|_{L^s(\mathbb{C})} \right], \end{aligned}$$

where we used Hölder's inequality (since $1/s = 1/2 + 1/\tilde{s}$) and estimate (4.1). The proof of (4.2) now follows from this last inequality and the fact that $\mu_z^2 - \mu_z^1 = \frac{1}{2}(\tau_+ - \tau_-)$. \square

Remark. We also have proved that

$$\begin{aligned} \sup_{z \in \mathbb{C}} \left\| \frac{\partial \mu_2(z, \cdot)}{\partial \bar{z}} - \frac{\partial \mu_1(z, \cdot)}{\partial \bar{z}} \right\|_{L^{\tilde{s}}(\mathbb{C})} &\leq c(D, N, s) \left[\left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^s(\mathbb{C})} \right. \\ &\quad \left. + \|h_2 - h_1\|_{L^s(\mathbb{C})} \right]. \end{aligned}$$

We will need the following consequence of Lemma 4.1.

Lemma 4.2. *Let v_1, v_2 be two potentials satisfying (1.3), (1.8), (1.10), with $\|v_j\|_{m,1} \leq N$. Let h_1, h_2 be the corresponding scattering amplitude and $\mu_1(z, \lambda), \mu_2(z, \lambda)$ the corresponding Faddeev functions. Let p, p' such that $1 < p < 2 < p' < \infty$, $1/p + 1/p' = 1$. Then*

$$(4.5) \quad \|\mu_2(\cdot, 0) - \mu_1(\cdot, 0)\|_{L^\infty(D)} \leq c(D, N, p) \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^p(\mathbb{C}) \cap L^{p'}(\mathbb{C})}.$$

Proof. We recall again that if $v \in W^{m,1}(\mathbb{R}^2)$, $m > 2$, with $\text{supp } v \subset D$ then $v \in L^p(D)$ for $p \in [1, \infty]$; in particular, from Proposition 2.1, this yields $h(\lambda)/\bar{\lambda} \in L^p(\mathbb{C})$, for $1 < p < \infty$.

We write, as in the preceding proof,

$$(4.6) \quad \nu(z, \lambda) = \mu_2(z, \lambda) - \mu_1(z, \lambda),$$

which satisfy the non-homogeneous $\bar{\partial}$ -equations (4.4). From this equation we obtain

$$\begin{aligned}
(4.7) \quad |\nu(z, 0)| &= \frac{1}{\pi} \left| \int_{\mathbb{C}} \frac{e_{-\lambda}(z)}{4\pi\lambda} \frac{h_1(\lambda)}{\bar{\lambda}} \overline{\nu(z, \lambda)} d\operatorname{Re}\lambda d\operatorname{Im}\lambda \right. \\
&\quad \left. + \int_{\mathbb{C}} \frac{e_{-\lambda}(z)}{4\pi\lambda} \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \overline{\mu_2(z, \lambda)} d\operatorname{Re}\lambda d\operatorname{Im}\lambda \right| \\
&\leq \frac{1}{4\pi^2} \sup_{z \in \mathbb{C}} \|\nu(z, \cdot)\|_{L^r} \left\| \frac{h_1(\lambda)}{\lambda\bar{\lambda}} \right\|_{L^{r'}} \\
&\quad + \frac{1}{4\pi^2} \sup_{z \in \mathbb{C}} \|\mu_2(z, \cdot)\|_{L^\infty} \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\lambda\bar{\lambda}} \right\|_{L^1}
\end{aligned}$$

where $1/r + 1/r' = 1$, $1 < r' < 2 < r < \infty$. The number $s = 2r/(r+2)$ can be chosen $s < 2$ and as close to 2 as wanted, by taking r big enough.

Then

$$(4.8) \quad \left\| \frac{h_1(\lambda)}{\lambda\bar{\lambda}} \right\|_{L^{r'}(|\lambda| < R)} \leq \left\| \frac{h_1(\lambda)}{\bar{\lambda}} \right\|_{L^p} \left\| \frac{1}{\lambda} \right\|_{L^q(|\lambda| < R)} \leq c(N, r),$$

where we have chosen $p > 2$ such that $\|h_1(\lambda)/\bar{\lambda}\|_{L^p} \leq c(N, p)$ from (2.10) and also, since $1/q = 1/r' - 1/p = 1 - 1/r - 1/p$, q can be chosen less than 2 by taking r big enough depending on p . With the same choice of p, q we also obtain

$$(4.9) \quad \left\| \frac{h_1(\lambda)}{\lambda\bar{\lambda}} \right\|_{L^{r'}(|\lambda| > R)} \leq \left\| \frac{h_1(\lambda)}{\bar{\lambda}} \right\|_{L^q} \left\| \frac{1}{\lambda} \right\|_{L^p(|\lambda| > R)} \leq c(N, r).$$

From Lemma 4.1 with $r = \tilde{s}$ we get

$$(4.10) \quad \sup_{z \in \mathbb{C}} \|\nu(z, \cdot)\|_{L^r} \leq c(D, N, r) \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^s(\mathbb{C})},$$

and from (2.11)

$$(4.11) \quad \sup_{z, \lambda \in \mathbb{C}} |\mu_2(z, \lambda)| \leq c(D, N).$$

Finally

$$\begin{aligned}
(4.12) \quad \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\lambda\bar{\lambda}} \right\|_{L^1} &\leq \left\| \frac{1}{\lambda} \right\|_{L^p(|\lambda| > R)} \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^{p'}} \\
&\quad + \left\| \frac{1}{\lambda} \right\|_{L^{p'}(|\lambda| < R)} \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^p},
\end{aligned}$$

by taking $p' = s$ and p such that $1/p + 1/p' = 1$. Now (4.5) follow from (4.6)–(4.12); this finishes the proof of Lemma 4.2. \square

5. PROOF OF THEOREMS 1.1 AND 1.2

Proof of Theorem 1.1. We begin with a remark, which take inspiration from Problem 1 at non-zero energy (see, for instance, [18]).

Let $v(z)$ be a potential which satisfies the hypothesis of Theorem 1.1 and $\mu(z, \lambda)$ the corresponding Faddeev functions. Since $\mu(z, \lambda)$ satisfies (2.11), the $\bar{\partial}$ -equation (2.8) and $h(\lambda)$ decreases at infinity like in Lemma 2.2, it is possible to write the following development:

$$(5.1) \quad \mu(z, \lambda) = 1 + \frac{\mu_{-1}(z)}{\lambda} + O\left(\frac{1}{|\lambda|^2}\right), \quad \lambda \rightarrow \infty,$$

for some function $\mu_{-1}(z)$. If we insert (5.1) into equation (2.1), for $\psi(z, \lambda) = e^{iz\lambda}\mu(z, \lambda)$, we obtain, letting $\lambda \rightarrow \infty$,

$$(5.2) \quad v(z) = 4i \frac{\partial \mu_{-1}(z)}{\partial \bar{z}}, \quad z \in \mathbb{C}.$$

We can write this in a more explicit form, using the following integral equation (a consequence of (2.8)):

$$\mu(z, \lambda) - 1 = \frac{1}{8\pi^2 i} \int_{\mathbb{C}} \frac{h(\lambda')}{(\lambda' - \lambda)\lambda'} e_{-\lambda'}(z) \overline{\mu(z, \lambda')} d\lambda' d\bar{\lambda}'.$$

By Lebesgue's dominated convergence (using (2.12)) we obtain

$$\mu_{-1}(z) = -\frac{1}{8\pi^2 i} \int_{\mathbb{C}} \frac{h(\lambda)}{\bar{\lambda}} e_{-\lambda}(z) \overline{\mu(z, \lambda)} d\lambda d\bar{\lambda},$$

and the explicit formula

$$(5.3) \quad v(z) = \frac{1}{2\pi^2} \int_{\mathbb{C}} e_{-\lambda}(z) \left(i h(\lambda) \overline{\mu(z, \lambda)} - \frac{h(\lambda)}{\lambda} \overline{\left(\frac{\partial \mu(z, \lambda)}{\partial z} \right)} \right) d\lambda d\bar{\lambda}.$$

Formula (5.3) for v_1 and v_2 yields

$$\begin{aligned} v_2(z) - v_1(z) = \frac{1}{2\pi^2} \int_{\mathbb{C}} e_{-\lambda}(z) & \left[i(h_2(\lambda) - h_1(\lambda)) \overline{\mu_2(z, \lambda)} \right. \\ & + i h_1(\lambda) (\overline{\mu_2(z, \lambda)} - \overline{\mu_1(z, \lambda)}) \\ & - \frac{h_2(\lambda) - h_1(\lambda)}{\lambda} \overline{\left(\frac{\partial \mu_2(z, \lambda)}{\partial z} \right)} \\ & \left. - \frac{h_1(\lambda)}{\bar{\lambda}} \overline{\left(\frac{\partial \mu_2(z, \lambda)}{\partial z} - \frac{\partial \mu_1(z, \lambda)}{\partial z} \right)} \right] d\lambda d\bar{\lambda}. \end{aligned}$$

Then, using several times Hölder's inequality, we find

$$\begin{aligned}
|v_2(z) - v_1(z)| &\leq \frac{1}{2\pi^2} \left(\|\mu_2(z, \cdot)\|_{L^\infty} \|h_2 - h_1\|_{L^1} \right. \\
&\quad + \|h_1\|_{L^{\tilde{p}'}} \|\mu_2(z, \cdot) - \mu_1(z, \cdot)\|_{L^{\tilde{p}}} \\
&\quad + \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^p} \left\| \frac{\partial \mu_2(z, \cdot)}{\partial z} \right\|_{L^{p'}} \\
&\quad \left. + \left\| \frac{h_1(\lambda)}{\bar{\lambda}} \right\|_{L^{\tilde{p}'}} \left\| \frac{\partial \mu_2(z, \cdot)}{\partial z} - \frac{\partial \mu_1(z, \cdot)}{\partial z} \right\|_{L^{\tilde{p}}} \right),
\end{aligned}$$

for $1 < p < 2$, \tilde{p} defined as in (2.3) and $1/p + 1/p' = 1/\tilde{p} + 1/\tilde{p}' = 1$. From (2.11), (2.10), the continuity of h_j and Lemma 2.2, [19, Lemma 2.1] (see the end of the proof of Lemma 4.1 for more details), Lemma 4.1, Propositions 3.4 and 3.3 we finally obtain

$$\begin{aligned}
\|v_2 - v_1\|_{L^\infty(D)} &\leq c(D, N, m, p) \left(\log(3 + \|\Phi_2 - \Phi_1\|_{H^{1/2} \rightarrow H^{-1/2}}^{-1})^{-(m-2)} \right. \\
&\quad + \log(3 + \|\Phi_2 - \Phi_1\|_{H^{1/2} \rightarrow H^{-1/2}}^{-1})^{-(m+1-2/p)} \\
&\quad \left. + \log(3 + \|\Phi_2 - \Phi_1\|_{H^{1/2} \rightarrow H^{-1/2}}^{-1})^{-(m-2/p)} \right) \\
&\leq c(D, N, m, p) \log(3 + \|\Phi_2 - \Phi_1\|_{H^{1/2} \rightarrow H^{-1/2}}^{-1})^{-(m-2)}.
\end{aligned}$$

This finishes the proof of Theorem 1.1. \square

Proof of Theorem 1.2. We first extend σ on the whole plane by putting $\sigma(x) = 1$ for $x \in \mathbb{R}^2 \setminus D$ (this extension is smooth by our hypothesis on σ). Now since $\sigma_j|_{\partial D} = 1$ and $\frac{\partial \sigma_j}{\partial \nu}|_{\partial D} = 0$ for $j = 1, 2$, from (1.5) we deduce that

$$(5.4) \quad \Phi_j = \Lambda_j, \quad j = 1, 2.$$

In addition, from (2.13) we get

$$(5.5) \quad \lim_{\lambda \rightarrow 0} \mu_j(z, \lambda) = \sigma_j^{1/2}(z), \quad j = 1, 2;$$

thus we obtain, using the fact that σ_j is bounded from above and below, for $j = 1, 2$,

$$\begin{aligned}
(5.6) \quad \|\sigma_2 - \sigma_1\|_{L^\infty(D)} &\leq c(N) \|\sigma_2^{1/2} - \sigma_1^{1/2}\|_{L^\infty(D)} \\
&= c(N) \|\mu_2(\cdot, 0) - \mu_1(\cdot, 0)\|_{L^\infty(D)}.
\end{aligned}$$

Now fix $\alpha < m$ and take p such that

$$\max \left(1, \frac{2}{m - \alpha + 1} \right) < p < 2.$$

From Lemma 4.2 we have

$$(5.7) \quad \|\mu_2(\cdot, 0) - \mu_1(\cdot, 0)\|_{L^\infty(D)} \leq c(D, N, p) \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^p(\mathbb{C}) \cap L^{p'}(\mathbb{C})},$$

where $1/p + 1/p' = 1$. From Proposition 3.3

$$\begin{aligned} \left\| \frac{h_2(\lambda) - h_1(\lambda)}{\bar{\lambda}} \right\|_{L^p(\mathbb{C}) \cap L^{p'}(\mathbb{C})} &\leq c(D, N, p) \log(3 + \|\Phi_2 - \Phi_1\|_{H^{1/2} \rightarrow H^{-1/2}}^{-1})^{-(m+1-2/p)} \\ &\leq c(D, N, p) \log(3 + \|\Phi_2 - \Phi_1\|_{H^{1/2} \rightarrow H^{-1/2}}^{-1})^{-\alpha} \\ &= c(D, N, p) \log(3 + \|\Lambda_2 - \Lambda_1\|_{H^{1/2} \rightarrow H^{-1/2}}^{-1})^{-\alpha}, \end{aligned}$$

from (5.4) and since $\alpha < m + 1 - \frac{2}{p}$. Theorem 1.2 is thus proved. \square

REFERENCES

- [1] G. Alessandrini, *Stable determination of conductivity by boundary measurements*, Appl. Anal. **27**, 1988, no. 1, 153–172.
- [2] K. Astala, L. Päivärinta, *Calderón’s inverse conductivity problem in the plane*, Ann. Math. **163**, 2006, 265–299.
- [3] J. A. Barceló, T. Barceló, A. Ruiz, *Stability of the inverse conductivity problem in the plane for less regular conductivities*, J. Diff. Equations **173**, 2001, 231–270.
- [4] T. Barceló, D. Faraco, A. Ruiz, *Stability of Calderón inverse conductivity problem in the plane*, J Math Pures Appl. **88**, 2007, no. 6, 522–556.
- [5] R. Beals, R. R. Coifman, *Multidimensional inverse scatterings and nonlinear partial differential equations*, Pseudodifferential operators and applications (Notre Dame, Ind., 1984), 45–70, Proc. Sympos. Pure Math., **43**, Amer. Math. Soc., Providence, RI, 1985.
- [6] A. L. Bukhgeim, *Recovering a potential from Cauchy data in the two-dimensional case*, J. Inverse Ill-Posed Probl. **16**, 2008, no. 1, 19–33.
- [7] A. P. Calderón, *On an inverse boundary problem*, Seminar on Numerical Analysis and its Applications to Continuum Physics, Soc. Brasileira de Matematica, Rio de Janeiro, 1980, 61–73.
- [8] B. A. Dubrovin, I. M. Krichever, S. P. Novikov, *The Schrödinger equation in a periodic field and Riemann surfaces*, Dokl. Akad. Nauk SSSR **229**, 1976, no. 1, 15–18.
- [9] L. D. Faddeev, *Growing solutions of the Schrödinger equation*, Dokl. Akad. Nauk SSSR **165**, 1965, no. 3, 514–517.
- [10] I. M. Gel’fand, *Some aspects of functional analysis and algebra*, Proceedings of the International Congress of Mathematicians, Amsterdam, 1954, **1**, 253–276. Erven P. Noordhoff N.V., Groningen; North-Holland Publishing Co., Amsterdam.
- [11] P. G. Grinevich, S. P. Novikov, *Two-dimensional “inverse scattering problem” for negative energies and generalized-analytic functions. I. Energies below the ground state*, Funct. Anal. and Appl. **22**, 1988, no. 1, 19–27.
- [12] G. M. Henkin, R. G. Novikov, *The $\bar{\partial}$ -equation in the multidimensional inverse scattering problem*, Russian Mathematical Surveys **42**, 1987, no. 3, 109–180.
- [13] M. Isaev, *Exponential instability in the Gel’fand inverse problem on the energy intervals*, J. Inverse Ill-Posed Probl. **19**, 2011, no. 3, 453–472; e-print arXiv:1012.2193.

- [14] L. Liu, *Stability Estimates for the Two-Dimensional Inverse Conductivity Problem*, Ph.D. thesis, Department of Mathematics, University of Rochester, New York, 1997.
- [15] N. Mandache, *Exponential instability in an inverse problem of the Schrödinger equation*, Inverse Problems **17**, 2001, no. 5, 1435–1444.
- [16] A. Nachman, *Global uniqueness for a two-dimensional inverse boundary value problem*, Ann. Math. **143**, 1996, 71–96.
- [17] R. G. Novikov, *Multidimensional inverse spectral problem for the equation $-\Delta\psi + (v(x) - Eu(x))\psi = 0$* , Funkt. Anal. i Pril. **22**, 1988, no. 4, 11–22 (in Russian); English Transl.: Funct. Anal. and Appl. **22**, 1988, no. 4, 263–272.
- [18] R. G. Novikov, *The inverse scattering problem on a fixed energy level for the two-dimensional Schrödinger operator*, J. Funct. Anal. **103**, 1992, no. 2, 409–463.
- [19] R. G., Novikov, *Approximate solution of the inverse problem of quantum scattering theory with fixed energy in dimension 2*, (Russian) Tr. Mat. Inst. Steklova **225**, 1999, Solitony Geom. Topol. na Perekrest., 301–318; translation in Proc. Steklov Inst. Math. **225**, 1999, no. 2, 285–302.
- [20] R. G. Novikov, *Formulae and equations for finding scattering data from the Dirichlet-to-Neumann map with nonzero background potential*, Inv. Problems 21, 2005, no. 1, 257–270.
- [21] R. G. Novikov, *New global stability estimates for the Gel’fand-Calderon inverse problem*, Inv. Problems **27**, 2011, no. 1, 015001.
- [22] R. G. Novikov, N. N. Novikova, *On stable determination of potential by boundary measurements*, ESAIM: Proc. **26**, 2009, 94–99.
- [23] R. G. Novikov, M. Santacesaria, *A global stability estimate for the Gel’fand-Calderón inverse problem in two dimensions*, J. Inverse Ill-Posed Probl. **18**, 2010, no. 7, 765–785.
- [24] I. N. Vekua, *Generalized Analytic Functions*, Pergamon Press Ltd. 1962.

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